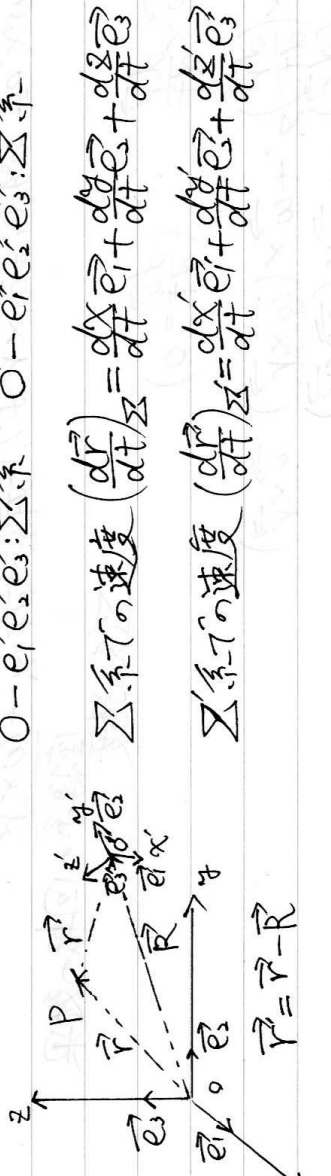


V 座標系の変換

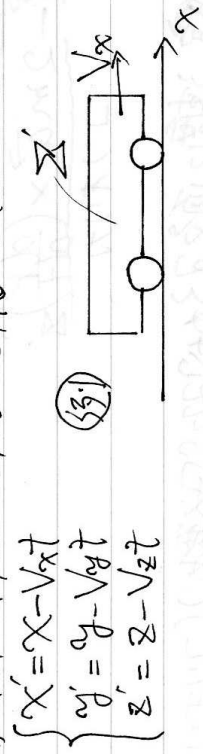


(1) カリル変換  $\vec{R} = \vec{R}_0 + \vec{v}t$  のとき,  $\vec{r}' = \vec{r} - \vec{R}$  を t の回数を消すと,

$$\left(\frac{d\vec{r}'}{dt}\right)_{\Sigma'} = \left(\frac{d\vec{r}}{dt}\right)_{\Sigma}$$

Σ系の運動方程式  $m \left(\frac{d\vec{r}'}{dt}\right)_{\Sigma'} = \vec{F}$  のとき, Σ系では  $m \left(\frac{d\vec{r}}{dt}\right)_{\Sigma} = \vec{F} \Rightarrow$  運動方程式は変換して

特に,  $\vec{e}_1' = \vec{e}_1, \vec{e}_2' = \vec{e}_2, \vec{e}_3' = \vec{e}_3, \vec{R}_0 = 0$  といわれる座標系の変換をガリリ変換という



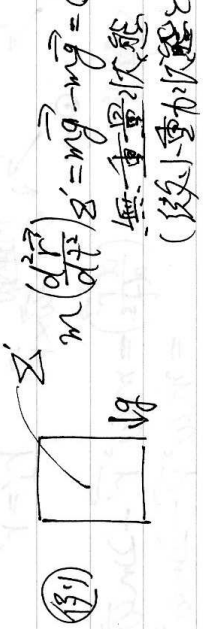
$$\begin{cases} x' = x - vt \\ y' = y \\ z' = z - vt \end{cases}$$

Σ系で真上へ投げ上げた物体の運動  
 $x' = 0, y' = v_0 t - \frac{1}{2}gt^2$

ΣΣ

Σ系では X 方向の初速度  $v_0$  を持った物体が投げ上り

$$x' = x - vt, y' = y \quad \therefore x = vt, y = v_0 t - \frac{1}{2}gt^2$$

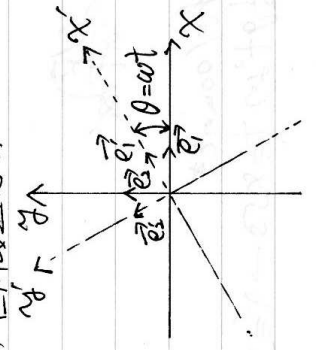


(2) 加速度系 -  $\vec{a} = \frac{d\vec{v}}{dt}$

$$\vec{r}' = \vec{r} - \vec{a}t^2, \left(\frac{d\vec{r}'}{dt}\right)_{\Sigma'} = \left(\frac{d\vec{r}}{dt}\right)_{\Sigma} - \vec{a}$$

$$\Sigma' \text{系での運動方程式 } m \left(\frac{d\vec{r}'}{dt}\right)_{\Sigma'} = \vec{F} - m\vec{a}$$

(3) 回転座標系



$$\vec{e}_1' = \cos\omega t \vec{e}_1 + \sin\omega t \vec{e}_2$$

$$\vec{e}_2' = -\sin\omega t \vec{e}_1 + \cos\omega t \vec{e}_2$$

$$\vec{e}_3' = \vec{e}_3$$

$$\frac{d\vec{e}_1'}{dt} = -\omega \sin\omega t \vec{e}_1 + \omega \cos\omega t \vec{e}_2 = \omega \vec{e}_3$$

$$\frac{d\vec{e}_2'}{dt} = -\omega \cos\omega t \vec{e}_1 - \omega \sin\omega t \vec{e}_2 = -\omega \vec{e}_3$$

$$\vec{r}' = x' \vec{e}_1' + y' \vec{e}_2' + z' \vec{e}_3'$$

$$\frac{d\vec{r}'}{dt} = \frac{dx'}{dt} \vec{e}_1' + \frac{dy'}{dt} \vec{e}_2' + \frac{dz'}{dt} \vec{e}_3' + x' \frac{d\vec{e}_1'}{dt} + y' \frac{d\vec{e}_2'}{dt} + z' \frac{d\vec{e}_3'}{dt}$$

$$= \left(\frac{d\vec{r}'}{dt}\right)_{\Sigma'} + \omega x' \vec{e}_3' - \omega y' \vec{e}_1'$$

(2次元-2次元)